



## **Dynamical System on Three – Dimensional Almost f-Cosymplectic Manifolds**

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### **Abstract**

The study addressed the Dynamical System on three-Almost f-Cosymplectic Manifolds and its applications in Mechanics using differential manifold techniques. It has been found that triple  $(M^3, \varphi, \xi, \eta, g)$  is Hamiltonian mechanical system on 3- Almost f-Cosymplectic Manifolds  $(M^3, \varphi, \xi, \eta, g)$ .

**Key words;** Almost f-Cosymplectic Manifolds, A differential Manifold, Dynamical System Equations.

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### **1 Introduction**

The geometric study of dynamical systems is an important topic in contemporary mathematics due to its applications in many branches, such, Mechanics and Theoretical Physics. If  $M$  is a differentiable manifold that corresponds

to the configuration space, a dynamical system can be locally given by a system of ordinary differential equations.

Mathematicians in [1],[2],[3] and [4] were introduced some important works dealt with concerned issues and its applications in many Mathematical fields, these motivated the researcher to deep insight in to dynamical system on three Almost f- Cosymplectic Manifold and its applications.

In this paper, we studied Dynamical System on Three- Almost f-Cosymplectic Manifold and its implementations in Mathematical Physics fields.

The work has been organized as follow, in section 1, we considered historical background as paper base, section 2 dealt with the study preliminary idea, section 3 devoted to study Three - dimensional Almost f-Cosymplectic Manifolds, section 4 devoted to study Lagrangian Dynamical System Equations on Three- Almost f-Cosymplectic Manifolds and section 5 is devoted to study Humiliation Dynamical System Equations Three- Almost f-Cosymplectic Manifolds.

## 2. Preliminary

In this preliminary chapter, we recalled basic definitions, results and formulas that we used them in the subsequent chapters of the paper.

### Definition Almost f-Cosymplectic Manifolds 2.1

Let  $M^n$  be a  $(2n + 1)$  –dimensional differentiable manifold equipped with a triple  $(\phi, \xi, \eta)$ , where  $\phi$  is a type of  $(1, 1)$  tensor field,  $\xi$  is a vector field,  $\eta$  is a 1-form on  $M$  such that

$$\eta(\xi) = 1, \quad \phi^2 = -1 + \eta \otimes \xi \quad (1)$$

which implies

$$\phi \xi = 0, \quad \eta \circ \phi = 0, \quad \text{rank}(\phi) = n - 1 \quad (2)$$

where  $\phi$  is a  $(1-1)$  tensor field  $\eta$  is a 1- form and the Riemannian metric  $g$ . It well known that

$$g(\phi X, \phi Y) = g(X, Y) - \eta(X)\eta(Y), \quad g(X, \xi) = \eta(X) \quad (3)$$

The fundamental 2- form of metric manifolds is defined by

$$g(\phi X, Y) = -g(X, \phi Y) \quad (4)$$

for any vector fields  $X, Y$  on  $M$ .

i) An almost contact metric manifold  $(M, \phi, \xi, \eta, g)$  such that  $d\omega = -f\eta \wedge \omega$  and  $d\eta = 0$  for smooth function  $f$  satisfying  $df \wedge \eta = 0$ . If the almost  $f$ -cosymplectic structure on  $M$  is normal.

**Proposition 2.2**

An  $f$ -cosymplectic manifold is cosymplectic manifold if  $f$  vanishes along  $\xi$ .

**Lemma 2.3**

For a three-dimensional  $f$ -cosymplectic manifold  $M^3$ , we have

$$QY = \left(-3\tilde{f} - \frac{R}{2}\right)\eta(Y)\xi + \left(\tilde{f} - \frac{R}{2}\right)Y \quad (6)$$

Where  $R$  is the scalar curvature of  $M$ .

**Lemma 2.4** If  $f, g$  and  $k$ -form be respectively then

- (i)  $f \wedge g = -g \wedge f$
- (ii)  $(f \wedge g)(x) = f(x)g - g(x)f$
- (iii)  $(dx^i \wedge dx^j) \left(\frac{\partial}{\partial x^k}\right) = \frac{\partial x^i}{\partial x^k} dx^j - \frac{\partial x^j}{\partial x^k} dx^i$
- (iv)  $\left(\frac{\partial x^j}{\partial x^i}\right) = \delta_j^i = \begin{cases} 1 & , i = j \\ 0 & , i \neq j \end{cases}$

### 3. Three – Dimensional Almost f-Cosymplectic Manifolds

**Definition 3.1**

Let three-dimensional manifold  $M = \{(x, y, z) \in R^3, z \neq 0\}$ ; where  $(x, y, z)$  are the standard coordinates in  $R^3$ : The vector fields

$$e_1 = e^{z^2} \frac{\partial}{\partial x} = e^{\theta(z)} \frac{\partial}{\partial x}, \quad e_2 = e^{z^2} \frac{\partial}{\partial y} = e^{\theta(z)} \frac{\partial}{\partial y}, \quad e_3 = \frac{\partial}{\partial z} \quad (6)$$

Where  $e^{\theta(z)}$  is a smooth function on  $M$ .

$\{e_1, e_2, e_3\}$  linearly independent at each point of  $M$ : Let  $g$  be the Riemannian metric defined by

$$g(e_1, e_3) = g(e_2, e_3) = g(e_1, e_2) = 0 \\ g(e_1, e_1) = g(e_2, e_2) = g(e_3, e_3) = 1$$

and given by the tensor product

$$g = \frac{1}{e^{z^2}} (dx \otimes dx + dy \otimes dy) + dz \otimes dz \quad (7)$$

**Proposition 3.2**

Let  $\eta$  be the 1-form defined by  $\eta(X) = g(X, e_3)$  for any  $X \in \mathfrak{X}(M)$ .

Let  $\phi$  be the (1,1) tensor field defined by

$$\phi(e_1) = e_2, \quad \phi(e_2) = -e_1 \text{ and } \phi(e_3) = 0 \quad (8)$$

Then using the linearity of  $\phi$  and  $g$  we have

$$\eta(e_3) = 1; \quad \phi^2(X) = -X + \eta(X)e_3; \\ g(\phi Z, \phi W) = g(X; Y) - \eta(X)\eta(Y);$$

for any  $Z, W \in \phi(M)$ : Thus for  $e_3 = \xi$ ,  $(\phi, \xi, \eta, g)$  defines an almost contact metric structure on  $M$ : Now, by direct computations we obtain

$$[e_1, e_2] = 0; \quad [e_1, e_3] = -2Ze_1, \quad [e_2, e_3] = -2Ze_1 \quad (9)$$

Consequently, the Nijenhuis torsion of  $\phi$  is zero, i.e.,  $M$  is normal.

$$\omega = \frac{1}{e^{\theta(z)}} dx \wedge dy = \frac{1}{e^{z^2}} dx \wedge dy \quad (10)$$

$$d\omega = 2\dot{\theta}(z) \omega \wedge \eta$$

Therefore  $M$  is an  $f$ -cosymplectic manifold with  $f(x, y, z) = \dot{\theta}(z)$ .

### Proposition 3.3

the following expressions are given

$$\phi\left(e^{z^2} \frac{\partial}{\partial x}\right) = e^{z^2} \frac{\partial}{\partial y}, \quad \phi\left(e^{z^2} \frac{\partial}{\partial y}\right) = -e^{z^2} \frac{\partial}{\partial x}, \quad \phi\left(\frac{\partial}{\partial z}\right) = 0 \quad (11)$$

The dual form  $\phi^*$  of the above  $\phi$  is as follows

$$\phi^*(e^{z^2} dx) = e^{z^2} dy, \quad \phi^*(e^{z^2} dy) = -e^{z^2} dx, \quad \phi^*(dz) = 0 \quad (12)$$

### Proposition 3.4

The vector fields

$$e_1 = \frac{\partial}{\partial x}, \quad e_2 = \frac{\partial}{\partial y} \text{ and } e_3 = \frac{\partial}{\partial z}$$

If  $\phi$  is defined a complex manifold  $\mathcal{M}$  then  $\phi^2 = \phi \circ \phi = -1$  or  $0$

**Proof:**

$$\phi^2(e_1) = \phi(\phi(e_1)) = \phi(e_2) = \phi(e_2) = -e_1 = -1$$

$$\phi^2(e_2) = \phi(\phi(e_2)) = \phi(-e_1) = -\phi(e_1) = -e_2 = -1$$

$$\phi^2(e_3) = \phi(\phi(e_3)) = \phi(0) = \phi(0) = 0$$

As can  $\phi^2$  is  $-1$  (complex) or  $0$

## 4. Lagrangian Dynamical System Equations on Three- Almost $f$ -Cosymplectic Manifolds

In this section we introduce the concept of Lagrangian Dynamical Systems. We start by the following definition.

### Definition 4.1

A Lagrangian function vector field  $X$  on  $\mathcal{M}$  is a smooth function

$L : T\mathcal{M} \rightarrow \mathbb{R}$  such that

$$i_X \phi_L = dE_L \quad (13)$$

### Definition 4.2

Let  $\xi$  be the vector field by

$$\xi = x \frac{\partial}{\partial x} + y \frac{\partial}{\partial y} + z \frac{\partial}{\partial z}, \quad x = \dot{x}, \quad y = \dot{y}, \quad z = \dot{z} \quad (14)$$

By Liouville vector field

space form  $(M^3, \phi, \xi, \eta, g)$ , we call the vector field determined by  $V = \phi\xi$  and calculated by

$$\begin{aligned} \phi\xi &= \phi\left(x \frac{\partial}{\partial x} + y \frac{\partial}{\partial y} + z \frac{\partial}{\partial z}\right) = X\phi\left(\frac{\partial}{\partial x}\right) + Y\phi\left(\frac{\partial}{\partial y}\right) + Z\phi\left(\frac{\partial}{\partial z}\right) \\ \phi\xi &= x \frac{\partial}{\partial y} - Y \frac{\partial}{\partial x} + Z(0) = Xe^{z^2} \frac{\partial}{\partial y} - Ye^{z^2} \frac{\partial}{\partial x} + Z(0) \\ \phi\xi &= Xe^{z^2} \frac{\partial}{\partial y} - Ye^{z^2} \frac{\partial}{\partial x} \end{aligned} \quad (15)$$

### Definition 4.3

1- Denote T by the **kinetic energy** and P by the **potential energy** of mechanics system on 3- Almost f-Cosymplectic Manifolds. Then we write by

$$L = T - P$$

2- Lagrangian function and by

$$E_L = V(L) - L \quad (16)$$

is called the interior product with  $\phi$ , or sometimes the insertion operator, or contraction by  $\phi$ .

**Definition 4.4**

The exterior vertical derivation  $d_\phi$  is defined by

$$d_\phi = [i_\phi, d] = i_\phi d - di_\phi \quad (17)$$

where  $d$  is the usual exterior derivation. For almost product structure  $\phi$  determined by the closed 3- Almost f-Cosymplectic Manifolds form is the closed 2-form given by

$$\phi_L = -dd_\phi L \quad (18)$$

such that

$$d_\phi = Xe^{z^2} \frac{\partial}{\partial y} - Ye^{z^2} \frac{\partial}{\partial x}$$

$$d_\phi L = \left( Xe^{z^2} \frac{\partial}{\partial y} - Ye^{z^2} \frac{\partial}{\partial x} \right) L = Xe^{z^2} \frac{\partial L}{\partial y} - Ye^{z^2} \frac{\partial L}{\partial x}$$

Thus we get

$$\phi_L = -d(d_\phi L) = -d \left( Xe^{z^2} \frac{\partial L}{\partial y} - Ye^{z^2} \frac{\partial L}{\partial x} \right)$$

$$\begin{aligned} \phi_L = & -Xe^{z^2} \frac{\partial^2 L}{\partial x \partial y} dx \wedge dx - Xe^{z^2} \frac{\partial^2 L}{\partial y \partial y} dy \wedge dy + Ye^{z^2} \frac{\partial^2 L}{\partial x \partial x} dx \wedge dy + Ye^{z^2} \frac{\partial^2 L}{\partial y \partial y} dy \wedge dy + \\ & Ze^{z^2} \frac{\partial^2 L}{\partial x \partial z} dx \wedge dz + Ze^{z^2} \frac{\partial^2 L}{\partial y \partial z} dy \wedge dz \end{aligned} \quad (19)$$

Because of the closed 3- Almost f-Cosymplectic Manifolds form  $\phi_L$  on 3- Almost f-Cosymplectic Manifolds space form  $(M^3, \phi, \xi, \eta, g)$  is para-symplectic structure, one may obtain

$$E_L = Xe^{z^2} \frac{\partial L}{\partial Y} - Ye^{z^2} \frac{\partial L}{\partial x} - L \quad (20)$$

Considering (0,1) we calculate

$$dE_L = d \left( Xe^{z^2} \frac{\partial L}{\partial Y} - Ye^{z^2} \frac{\partial L}{\partial x} - L \right)$$

$$\begin{aligned} dE_L = & -Xe^{z^2} \frac{\partial^2 L}{\partial x \partial y} dx - ye^{z^2} \frac{\partial^2 L}{\partial y \partial y} dy + Xe^{z^2} \frac{\partial^2 L}{\partial x \partial x} dx + Ye^{z^2} \frac{\partial^2 L}{\partial y \partial x} dy - \frac{\partial L}{\partial x} dx - \\ & Z \frac{\partial^2 L}{\partial x \partial z} dx - Ze^{z^2} \frac{\partial^2 L}{\partial y \partial z} dy - \frac{\partial L}{\partial y} dy \end{aligned} \quad (21)$$

Taking care of  $i_\xi \phi_L = dE_L$ , we have

$$\begin{aligned} & -Xe^{z^2} \frac{\partial^2 L}{\partial x \partial y} dx - ye^{z^2} \frac{\partial^2 L}{\partial y \partial y} dy - Ze^{z^2} \frac{\partial^2 L}{\partial y \partial z} dy + \frac{\partial L}{\partial x} dx + Xe^{z^2} \frac{\partial^2 L}{\partial x \partial x} dy + \\ & Ye^{z^2} \frac{\partial^2 L}{\partial y \partial x} dy + Ze^{z^2} \frac{\partial^2 L}{\partial x \partial z} dy + \frac{\partial L}{\partial y} dy = 0 \end{aligned} \quad (22)$$

$$\begin{aligned}
 & -\frac{\partial L}{\partial y} \left( X e^{z^2} \frac{\partial}{\partial x} dx + Y e^{z^2} \frac{\partial}{\partial y} dy + Z e^{z^2} \frac{\partial}{\partial z} dz \right) + \frac{\partial L}{\partial x} dx \\
 & + \frac{\partial L}{\partial x} \left( X e^{z^2} \frac{\partial}{\partial x} dy + Y e^{z^2} \frac{\partial}{\partial y} dy + Z e^{z^2} \frac{\partial}{\partial z} dy \right) + \frac{\partial L}{\partial y} dy = 0 \quad (23) \\
 & -\frac{\partial L}{\partial y} e^{z^2} \left( X \frac{\partial}{\partial x} + Y \frac{\partial}{\partial y} + Z \frac{\partial}{\partial z} \right) dx + \frac{\partial L}{\partial x} dx + \frac{\partial L}{\partial x} e^{z^2} \left( X \frac{\partial}{\partial x} + Y \frac{\partial}{\partial y} + Z \frac{\partial}{\partial z} \right) dy \\
 & + \frac{\partial L}{\partial y} dy = 0 \\
 & \left[ -e^{z^2} \frac{\partial L}{\partial y} \left( X \frac{\partial}{\partial x} + Y \frac{\partial}{\partial y} + Z \frac{\partial}{\partial z} \right) + \frac{\partial L}{\partial x} \right] dx + \left[ e^{z^2} \frac{\partial L}{\partial x} \left( X \frac{\partial}{\partial x} + Y \frac{\partial}{\partial y} + Z \frac{\partial}{\partial z} \right) + \frac{\partial L}{\partial y} \right] dy = 0
 \end{aligned}$$

If the curve  $\alpha : I \subset \mathbb{R} \rightarrow M^3$  be integral curve of  $\xi$ ,

$$\alpha = \frac{\partial}{\partial t} = X \frac{\partial}{\partial x} + Y \frac{\partial}{\partial y} + Z \frac{\partial}{\partial z} \quad (24)$$

which satisfies

$$\left[ -e^{z^2} \frac{\partial}{\partial t} \frac{\partial L}{\partial y} + \frac{\partial L}{\partial x} \right] dx + \left[ e^{z^2} \frac{\partial}{\partial t} \frac{\partial L}{\partial x} + \frac{\partial L}{\partial y} \right] dy = 0 \quad (25)$$

it follows equations

$$e^{z^2} \frac{\partial}{\partial t} \frac{\partial L}{\partial y} - \frac{\partial L}{\partial x} = 0, \quad e^{z^2} \frac{\partial}{\partial t} \frac{\partial L}{\partial x} + \frac{\partial L}{\partial y} = 0 \quad (26)$$

Finally one may say that the triple  $(M^3, \phi, \xi, \eta, g)$  is Lagrangian mechanical system on 3- Almost  $f$ -Cosymplectic Manifolds  $(M^3, \phi, \xi, \eta, g)$ .

## 5. Humiliation Dynamical System Equationsthree- Almost $f$ -Cosymplectic Manifolds

### Definition 5.1

The kinetic energy  $T = \frac{1}{2} m(\dot{x}^2 + \dot{y}^2 + \dot{z}^2)$  and  $P = m_i gh$  is Potential energy .then  $L = T + P$  is called the Hamiltonian function.

### Definition 5.2

Let  $X_H$  is unique vector field on  $H$  is Hamiltonian function and  $\phi^*$  is dual of  $\phi$  and  $\omega$  from on  $T^*M$  and

$$i_{X_H} = dH \quad (27)$$

Is called Hamiltonian dynamical equation .

Here, we present Hamiltonian equations on 3- Almost  $f$ -Cosymplectic Manifolds  $(M^3, \phi, \xi, \eta, g)$  such that

$$\omega = \frac{1}{2} (Xdx + Ydy + Zdz) \quad (28)$$

1-form on

Let  $\phi^*$  be an almost product structure defined by  $\lambda$  Liouville form determined by

$$\lambda = \phi^*(\omega) = \phi^* \left[ \frac{1}{2} (Xdx + Ydy + Zdz) \right] \quad (29)$$

$$\lambda = \phi^*(\omega) = \frac{1}{2} [X\phi^*(dx) + Y\phi^*(dy) + Z\phi^*(dz)]$$

$$\begin{aligned}\lambda &= \phi^*(\omega) = \frac{1}{2} [X(e^{z^2} dy) + Y(-e^{z^2} dx) + Z(0)] \\ \lambda &= \phi^*(\omega) = \frac{1}{2} [Xe^{z^2} dy - Ye^{z^2} dx]\end{aligned}\quad (30)$$

Differential of  $\lambda$

$$\varphi = -d\lambda = -d\left(\frac{1}{2} [Xe^{z^2} dy - Ye^{z^2} dx]\right)$$

It is known that if  $\varphi$  is a closed 2- form on  $T^*M^3$ , then  $\varphi_H$  is also a Symplectic structure on  $T^*M^3$

$$\varphi = -d\lambda = e^{z^2} dy \wedge dx \quad (31)$$

. Let  $(M^3, \phi, \xi, \eta, g)$  3- Almost  $f$ -Cosymplectic Manifolds form  $\varphi$ .

Suppose that Hamiltonian vector field  $X_H$  associated to Hamiltonian energy  $H$  is given by

$$X_H = X \frac{\partial}{\partial x} + Y \frac{\partial}{\partial y} + Z \frac{\partial}{\partial z} \quad (32)$$

Calculates a value  $X_H$  and  $\varphi$

$$\varphi(Z_H) = (e^{z^2} dy \wedge dx) \left( X \frac{\partial}{\partial x} + Y \frac{\partial}{\partial y} + Z \frac{\partial}{\partial z} \right)$$

$$\varphi(Z_H) = -e^{z^2} (dy \wedge dx) \left( x \frac{\partial}{\partial x} \right) + dy \wedge dx \left( y \frac{\partial}{\partial y} \right) + dy \wedge dx \left( z \frac{\partial}{\partial z} \right)$$

$$\varphi(Z_H) = -Xe^{z^2} \frac{\partial}{\partial y} + Ye^{z^2} \frac{\partial}{\partial x} - Ze^{z^2} \frac{\partial}{\partial x} + Xz \frac{\partial}{\partial x} + Ze^{z^2} \frac{\partial}{\partial x} - Ze^{z^2} \frac{\partial}{\partial y} \quad (33)$$

Otherwise, one may calculate the differential of Hamiltonian energy as follows:

$$dH = \frac{\partial H}{\partial x} dx + \frac{\partial H}{\partial y} dy + \frac{\partial H}{\partial z} dz \quad (34)$$

From (33) and (34) with respect to  $i_{X_H} \varphi = dH$ , we find Hamiltonian vector field on 3- Almost  $f$ -Cosymplectic Manifolds space be

$$X_H = \frac{\partial H}{\partial y} \frac{\partial}{\partial x} - \frac{\partial H}{\partial x} \frac{\partial}{\partial y} \quad (35)$$

Suppose that the curve

$$\alpha: I \subset \mathbb{R} \rightarrow \mathbb{R}^{2n}$$

be an integral curve of Hamiltonian vector field  $Z_H$ , i.e.,

$$Z_H(\alpha(t)) = \dot{\alpha}, \quad t \in I.$$

In the local coordinates we have

$$\alpha(t) = (x(t), y(t), z(t)),$$

$$\dot{\alpha}(t) = \frac{dx}{dt} \frac{\partial}{\partial x} + \frac{dy}{dt} \frac{\partial}{\partial y} + \frac{dz}{dt} \frac{\partial}{\partial z} \quad (36)$$

Now, by means of (26), from (27) and (28), we deduce the equations so-called Hamiltonian equations

$$\frac{dx}{dt} = e^{z^2} \frac{\partial H}{\partial y}, \quad \frac{dy}{dt} = -e^{z^2} \frac{\partial H}{\partial x} \quad (37)$$

## 6. Conclusion:

Finally, we concluded that the triple  $(M^3, \phi, \xi, \eta, g)$  is Hamiltonianmechanical system on 3- Almost  $f$ -Cosymplectic Manifolds  $(M^3, \phi, \xi, \eta, g)$  from the above deductions.

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